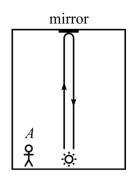
Time Dilation

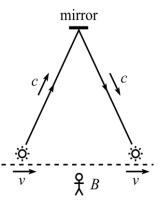


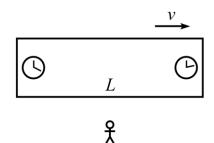
- Light moves at speed *c* in all frames
- Two events that occur simultaneously in one reference frame may not be simultaneous in another frame
- Two events that occur <u>at the same location</u> in frame *A*, separated by a time t_A , will take a greater time $t_B = \gamma t_A$ in frame *B* moving at velocity v with respect to frame *A*

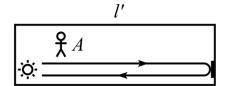
$$\gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}$$
 (with $\gamma \ge 1$)

Head Start Effect

- Two clocks are positioned at the end of a train of length *L* (as measured in its own frame. The clocks are synchronized in the train's reference frame
- The train travels past an observer at speed v. The observer will see the *rear clock* showing a higher reading than the front clock by $\Delta t = \frac{Lv}{c^2}$
- The rear clock does not tick faster than the front clock, it simply remains a fixed time ahead of the front clock







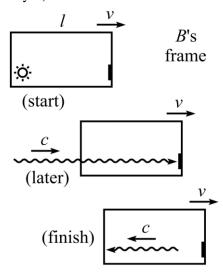
Length Contraction

- An object in frame A moves with velocity v relative to frame B. The object has length l' along this direction of motion in A's frame. The object has a *shorter* length $l = \frac{l'}{v}$ in B's frame
- Length contraction does not apply in the directions transverse to the direction of motion (set by *v*) between two reference frames

Example

A muon moves from a height h straight down towards the Earth with a large velocity v. It decays after time T in its own frame. Can the muon reach the ground?

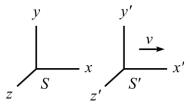
- In the Earth's frame, time dilation applies. The muon has to travel a distance *h* in time *γ T*
- In the muon's frame, length contraction kicks in. The muon must travel a distance $\frac{h}{\gamma}$ in time T



Lorentz Transformation

- Frame *S'* moves at velocity *v* relative to frame *S*
- The length and time between two events (an event is anything that has space and time coordinates) in *S*' are related to those in *S* by

$$\Delta x = \gamma (\Delta x' + v \Delta t') \qquad \Delta x' = \gamma (\Delta x - v \Delta t)$$
$$\Delta t = \gamma \left(\Delta t' + \frac{v}{c^2} \Delta x' \right) \qquad \Delta t' = \gamma \left(\Delta t - \frac{v}{c^2} \Delta x \right)$$



- Since there is no transverse length contraction, $\Delta y = \Delta y'$ and $\Delta z = \Delta z'$
- $\Delta x^2 c^2 \Delta t^2$ has the same value in all reference frames

Velocity Addition

• An object moves at speed v_1 with respect to frame S'. Frame S' moves at speed v_2 with respect to frame S in the same direction of motion as the object. The speed u of the object in frame S is

$$u = \frac{v_1 + v_2}{1 + \frac{v_1 v_2}{c^2}}$$

- An object moving at speed *c* in one reference frame moves at speed *c* in all reference frames
- An object moving slower than *c* in one reference frame moves slower than *c* in all reference frames

Momentum and Energy

• The relativistic energy and momentum of a particle moving with velocity \vec{v} are given by

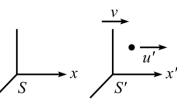
$$\vec{p} = \gamma m \vec{v}$$

$$E = \gamma m c^{2}$$
Slow speeds
$$\vec{p} \approx m \vec{v}$$

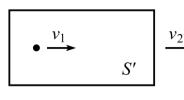
$$E \approx m c^{2} + \frac{1}{2} m v^{2}$$

• A particle moves at speed u' in frame S', which moves at speed v relative to frame S. Momentum $(\gamma_{u'}mu')$ and energy $(\gamma_{u'}mc^2)$ in S' is related to that in S by the same Lorentz transformation as space and time (with $\gamma \equiv \gamma_v$)

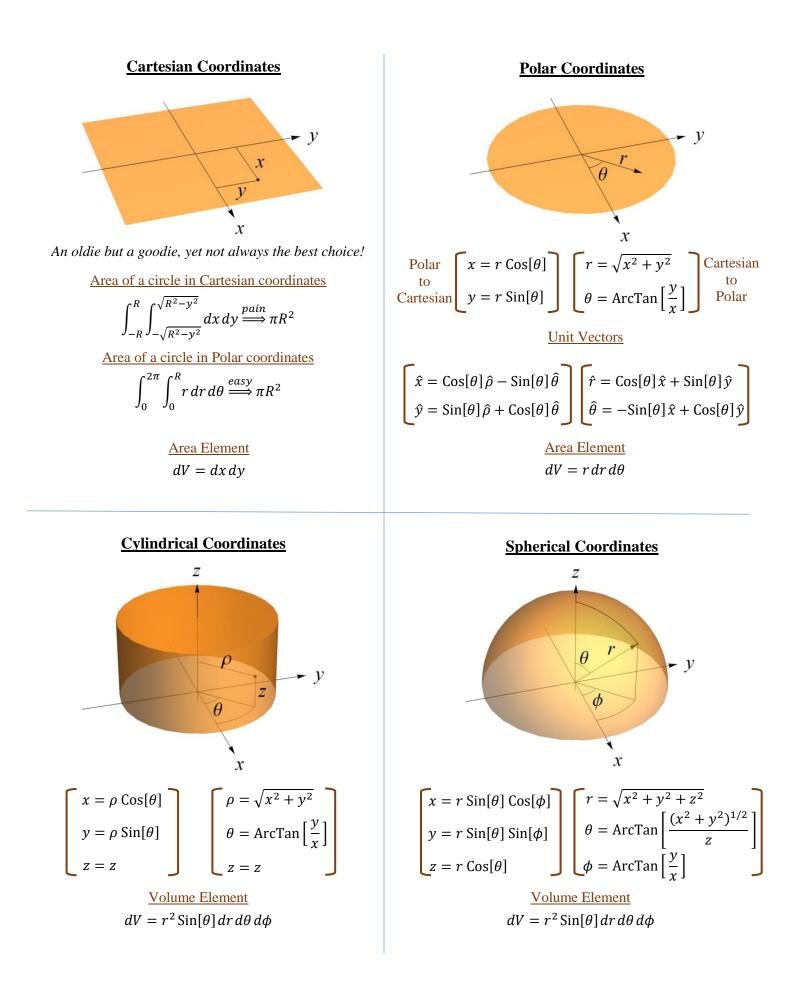
$$E = \gamma(E' + v p') \qquad E' = \gamma(E - v p) p = \gamma \left(p' + \frac{v}{c^2} E' \right) \qquad p' = \gamma \left(p - \frac{v}{c^2} E \right)$$

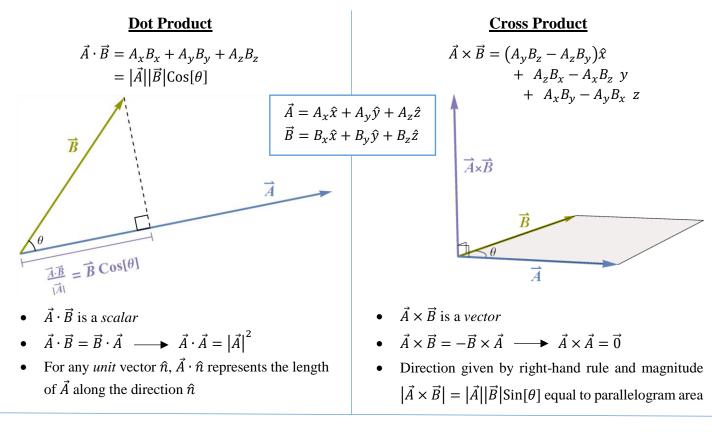


- Energy and momentum are conserved in collisions *within the same reference frame*
- The invariant mass formula $E^2 p^2 c^2 = m^2 c^4$ applies any mass *in any reference frame*

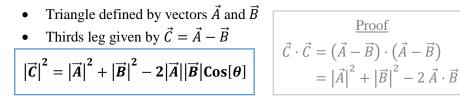


S

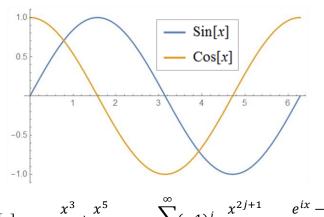




Law of Cosines



Trig Functions



 \vec{C}

À

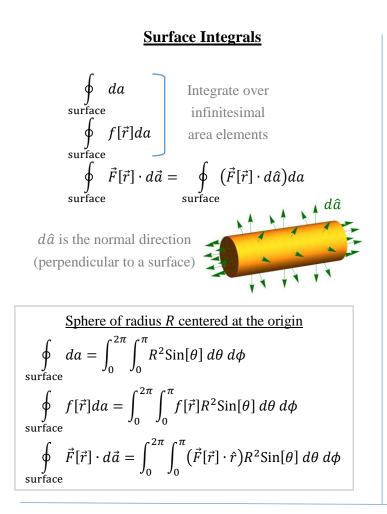
 \overrightarrow{B}

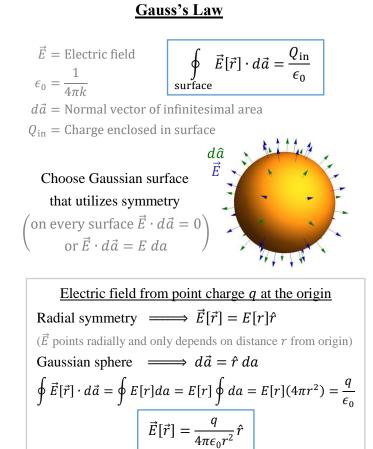
 θ

$$\operatorname{Sin}[x] = x - \frac{x^3}{3!} + \frac{x^5}{5!} - \dots = \sum_{j=0}^{\infty} (-1)^j \frac{x^{2j+1}}{(2j+1)!} = \frac{e^{ix} - e^{-ix}}{2}$$
$$\operatorname{Cos}[x] = 1 - \frac{x^2}{2!} + \frac{x^4}{4!} - \dots = \sum_{j=0}^{\infty} (-1)^j \frac{x^{2j}}{(2j)!} = \frac{e^{ix} + e^{-ix}}{2}$$

 $\operatorname{Cos}[0] = 1 \qquad \operatorname{Sin}[0] = 0$ $\operatorname{Cos}\left[\frac{\pi}{6}\right] = \frac{\sqrt{3}}{2} \qquad \operatorname{Sin}\left[\frac{\pi}{6}\right] = \frac{1}{2}$ $\operatorname{Cos}\left[\frac{\pi}{4}\right] = \frac{\sqrt{2}}{2} \qquad \operatorname{Sin}\left[\frac{\pi}{4}\right] = \frac{\sqrt{2}}{2}$ $\operatorname{Cos}\left[\frac{\pi}{3}\right] = \frac{1}{2} \qquad \operatorname{Sin}\left[\frac{\pi}{3}\right] = \frac{\sqrt{3}}{2}$ $\operatorname{Cos}\left[\frac{\pi}{2}\right] = 0 \qquad \operatorname{Sin}\left[\frac{\pi}{2}\right] = 1$ $\operatorname{Cos}[\pi - x] = -\operatorname{Cos}[x] \qquad \operatorname{Sin}[\pi - x] = \operatorname{Sin}[x]$ $\operatorname{Sin}[x + y] = \operatorname{Sin}[x]\operatorname{Cos}[y] + \operatorname{Cos}[x]\operatorname{Sin}[y]$ $\operatorname{Cos}[x + y] = \operatorname{Cos}[x]\operatorname{Cos}[y] - \operatorname{Sin}[x]\operatorname{Sin}[y]$ $\operatorname{Sin}[2x] = 2\operatorname{Sin}[x]\operatorname{Cos}[x]$ $\operatorname{Cos}[2x] = \operatorname{Cos}[x]^2 - \operatorname{Sin}[x]^2$

$$= 2\cos[x]^2 - 1 = 1 - 2\sin[x]^2$$



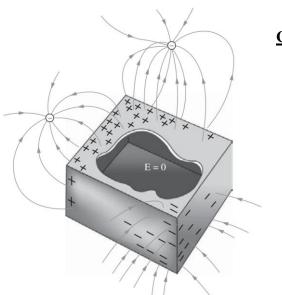


Example: Infinite Sheet of Charge

Infinite plane at x = 0 with uniform charge density σ Translational symmetry $\implies \vec{E}[\vec{r}] = E[x]\hat{x}$ (\vec{E} points away from plane and only depends on distance x from plane) Gaussian cube with side length 2s centered on the plane • $E[x]\hat{x} \cdot d\vec{a} = 0$ for all sides that intersect the plane • $E[x]\hat{x} \cdot d\vec{a} = E[x]da$ for both surfaces parallel to plane • $E[x]\hat{x} \cdot d\vec{a} = E[x]da$ for both surfaces parallel to plane • $E[x]\hat{x} \cdot d\vec{a} = 2E[x](4s^2) = \frac{\sigma(4s^2)}{\epsilon_0}$ $\vec{E}[\vec{r}] = \begin{cases} \frac{\sigma}{2\epsilon_0}\hat{x} & x > 0\\ -\frac{\sigma}{2\epsilon_0}\hat{x} & x < 0 \end{cases}$ Exactly on the plane, electric field must be zero by symmetry! Discontinuity in \vec{E} across a sheet of charge is always $\frac{\sigma}{\epsilon_0}$ Value of \vec{E} on the sheet is the average of its values on both sides

Electrostatics

Fundamental quantities: $\rho - \text{Charge distribution}$ $\vec{E} - \text{Force per unit charge on a test particle}$ $\phi - \frac{P \text{otential energy per unit charge}}{\text{to place a test particle}}$ Work required to assemble charge distribution: $U = \frac{1}{2} \sum_{i=1}^{N} \sum_{j \neq i} \frac{k q_i q_j}{r_{jk}} \quad \text{(Point charges)}$ $U = \frac{\epsilon_0}{2} \int E^2 dv = \frac{1}{2} \int \rho \phi dv \quad \text{(Charge distribution)} \quad \rho$



Conductors

Key properties

 \vec{E}

- $\vec{E} = \vec{0}$ inside a conductor
- $\rho = 0$ inside a conductor
- Any net charge resides on the surface
- A conductor surface is equipotential
- \vec{E} perpendicular to surface just outside a conductor

Example: Parallel Plates (Ignoring Edge Effects)

